

# ASTR 400/700: Stellar Astrophysics

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# Stellar Atmospheres

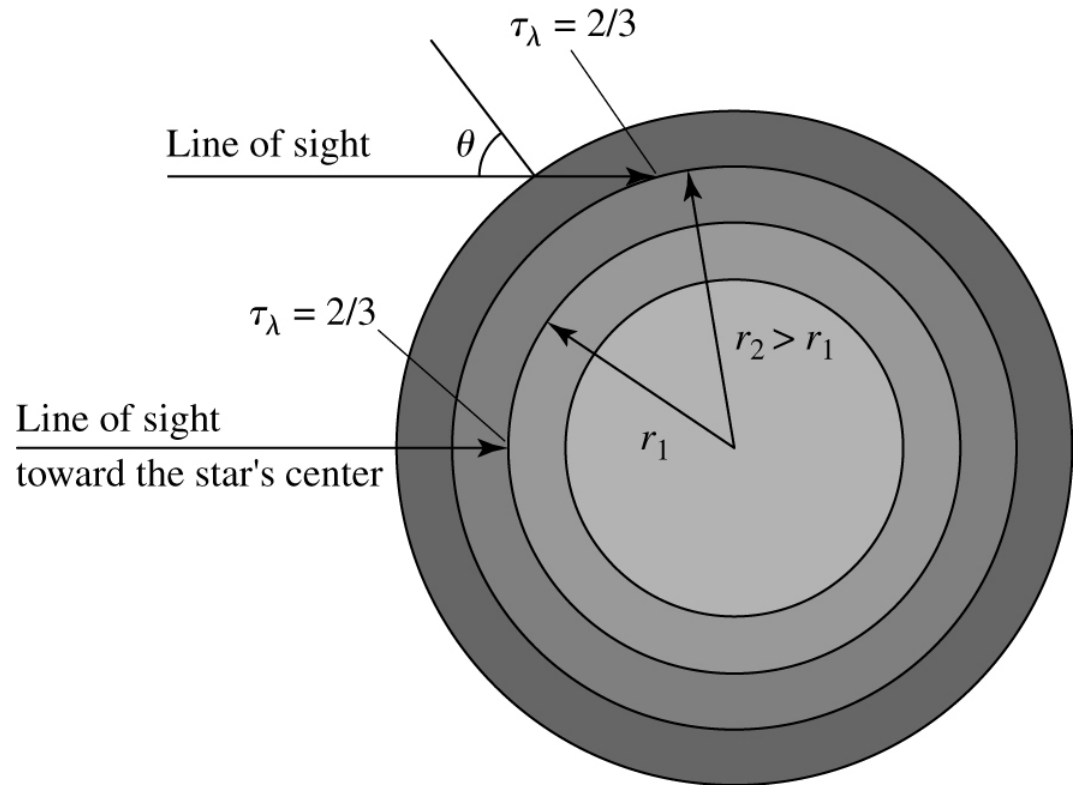
Chapter 9.4, 9.5

# Limb Darkening

- When looking at the center of the sun one can see “deeper” than when looking at the edge of the sun
  - Deeper is hotter
  - Hotter is brighter...



LIMB DARKENING



# The Transfer Equation

## Describes the passage of light through a Star's Atmosphere

- The Emission Coefficient  $j_\lambda$ 
  - Describes the rate of emission processes



$$dI_\lambda = j_\lambda \rho ds,$$

- Source of radiation  $S_\lambda = j_\lambda / \kappa_\lambda$ 
  - Ratio of rates of emission to absorption



$$dI_\lambda = -\kappa_\lambda \rho I_\lambda ds + j_\lambda \rho ds.$$
$$-\frac{1}{\kappa_\lambda \rho} \frac{dI_\lambda}{ds} = I_\lambda - \frac{j_\lambda}{\kappa_\lambda}.$$

- Equation of Radiative Transport



$$-\frac{1}{\kappa_\lambda \rho} \frac{dI_\lambda}{ds} = I_\lambda - S_\lambda.$$

# Plane-Parallel Atmosphere

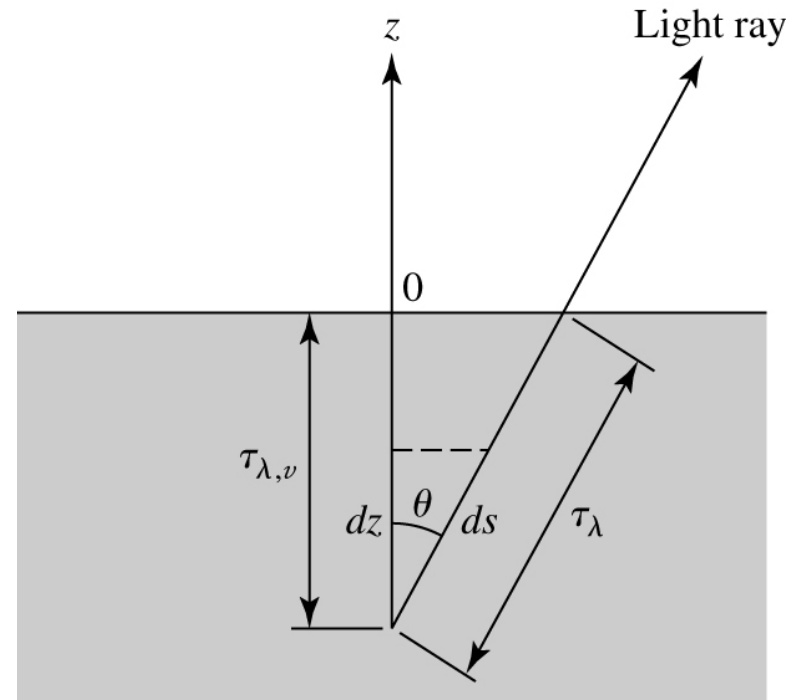
- Far from center of star the radius of curvature is large...consider plane-parallel (flat) slabs...
- Work in units of optical depth instead of distance...
- Z-axis is vertical  $z=0$  is at top of slab
  - Vertical Optical Depth

$$\tau_{\lambda,v}(z) \equiv \int_z^0 \kappa_{\lambda} \rho dz.$$

Further simplify by using mean opacity--->"Gray Atmosphere". No more dependence on wavelength.

In an equilibrium stellar atmosphere, every process of absorption is balanced by an inverse process of emission, **no net energy is subtracted from or added to the radiation field.**

$$\frac{dI_{\lambda}}{d\tau_{\lambda}} = I_{\lambda} - S_{\lambda}.$$

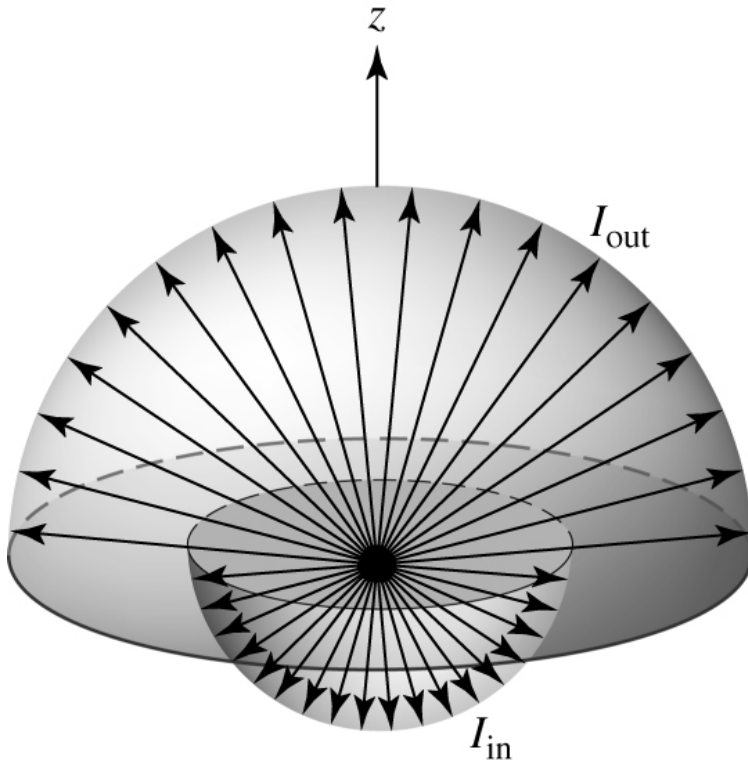


# Eddington Approximation

- Approximate  $I$  with  $I_{\text{out}}$  and  $I_{\text{in}}$
- Evaluate Temperature as a function of optical depth

$$\langle I \rangle = \frac{3\sigma}{4\pi} T_e^4 \left( \tau_v + \frac{2}{3} \right).$$

$$T^4 = \frac{3}{4} T_e^4 \left( \tau_v + \frac{2}{3} \right).$$



can be generalized to the statement that *when looking at a star, we see down to a vertical optical depth of  $\tau_v \approx 2/3$ , averaged over the disk of the star.* The importance of this for the formation and interpretation of spectral lines was discussed on page 254.

# Limb Darkening Revisited

$$\frac{dI_\lambda}{d\tau_\lambda} = I_\lambda - S_\lambda,$$

$$\frac{dI_\lambda}{d\tau_\lambda} e^{-\tau_\lambda} - I_\lambda e^{-\tau_\lambda} = -S_\lambda e^{-\tau_\lambda}$$

$$\frac{d}{d\tau_\lambda} (e^{-\tau_\lambda} I_\lambda) = -S_\lambda e^{-\tau_\lambda}$$

$$d(e^{-\tau_\lambda} I_\lambda) = -S_\lambda e^{-\tau_\lambda} d\tau_\lambda.$$

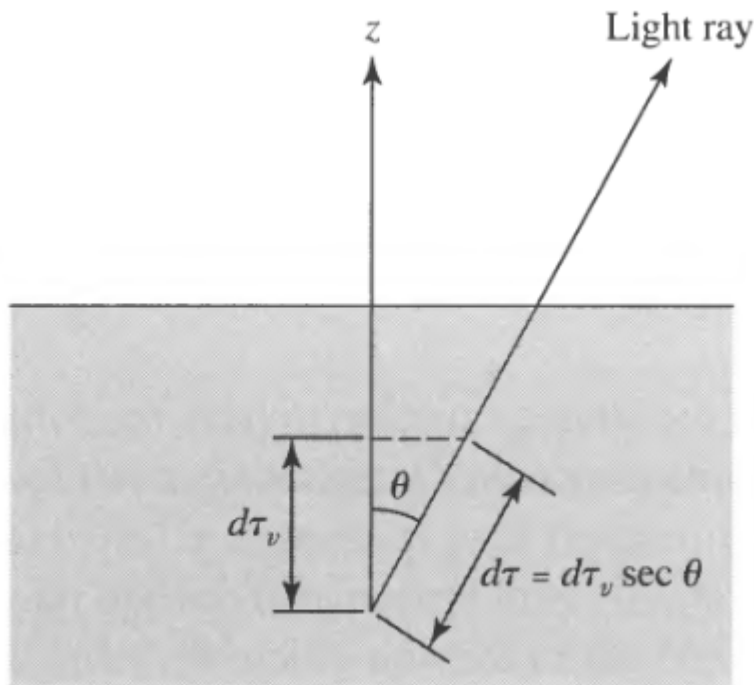
$$I_\lambda(0) = I_{\lambda,0} e^{-\tau_{\lambda,0}} - \int_{\tau_{\lambda,0}}^0 S_\lambda e^{-\tau_\lambda} d\tau_\lambda.$$

$$I(0) = I_0 e^{-\tau_{v,0} \sec \theta} - \int_{\tau_{v,0} \sec \theta}^0 S \sec \theta e^{-\tau_v \sec \theta} d\tau_v.$$

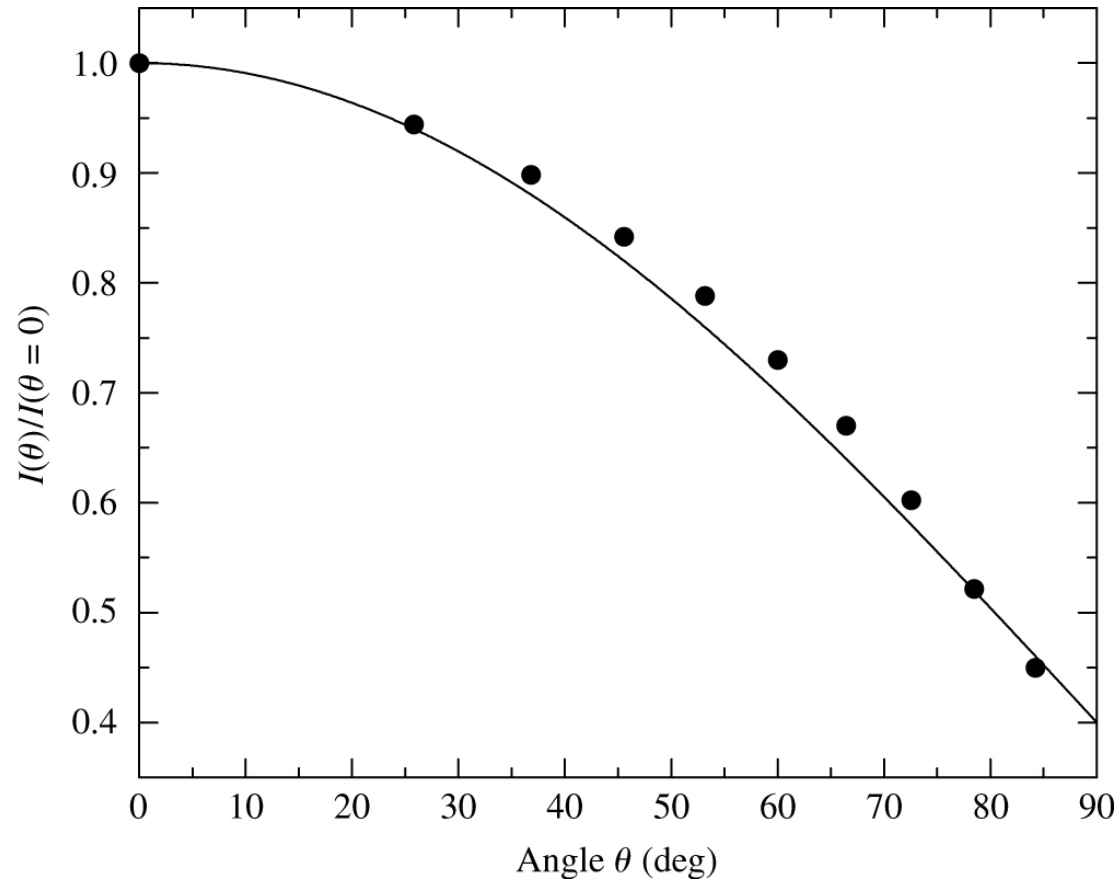
$$I(0) = \int_0^\infty S \sec \theta e^{-\tau_v \sec \theta} d\tau_v.$$

$$S = a + b\tau_v,$$

$$I_\lambda(0) = a_\lambda + b_\lambda \cos \theta,$$



# Eddington Approximation of Solar Limb Darkening





# Profiles of Spectral Lines

- Spectral Line conveys information about the environment in which it was formed

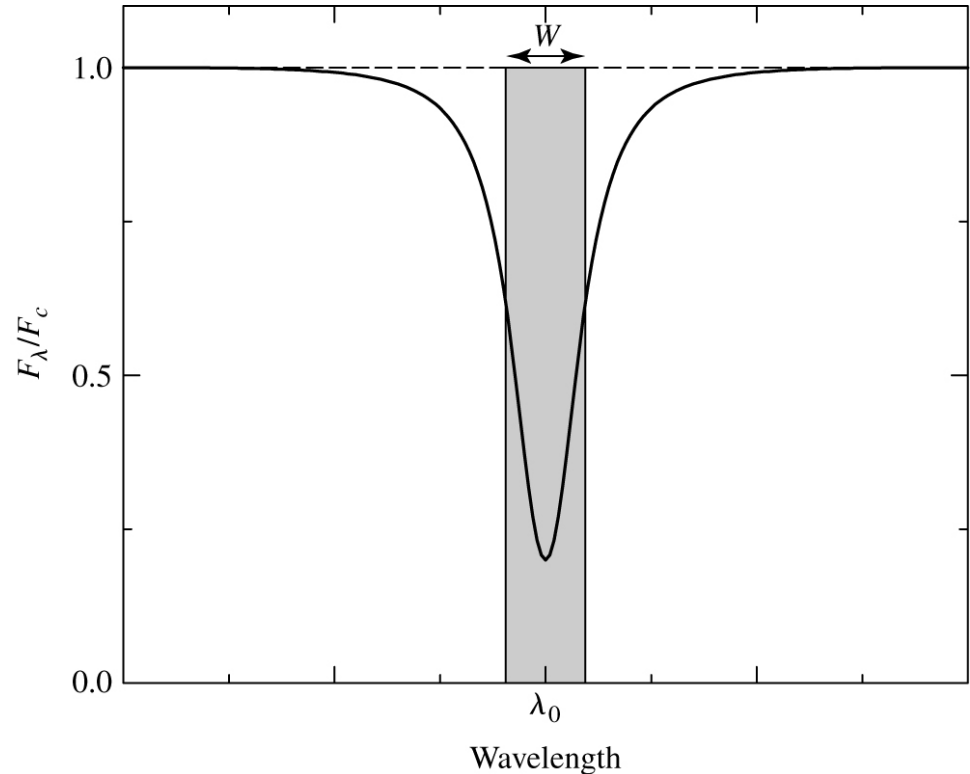
- Line width

$$W = \int \frac{F_c - F_\lambda}{F_c} d\lambda,$$

- Line Shape

– Broadening

– Profile



**Optically Thin Spectral Line.** Thus termed because there is no wavelength at which the radiant flux has been **completely** blocked.

# Sources of Broadening

- **Natural Broadening**

- From Uncertainty Principle
- $\sim 2-4 \times 10^{-5}$  nm

- **Doppler Broadening**

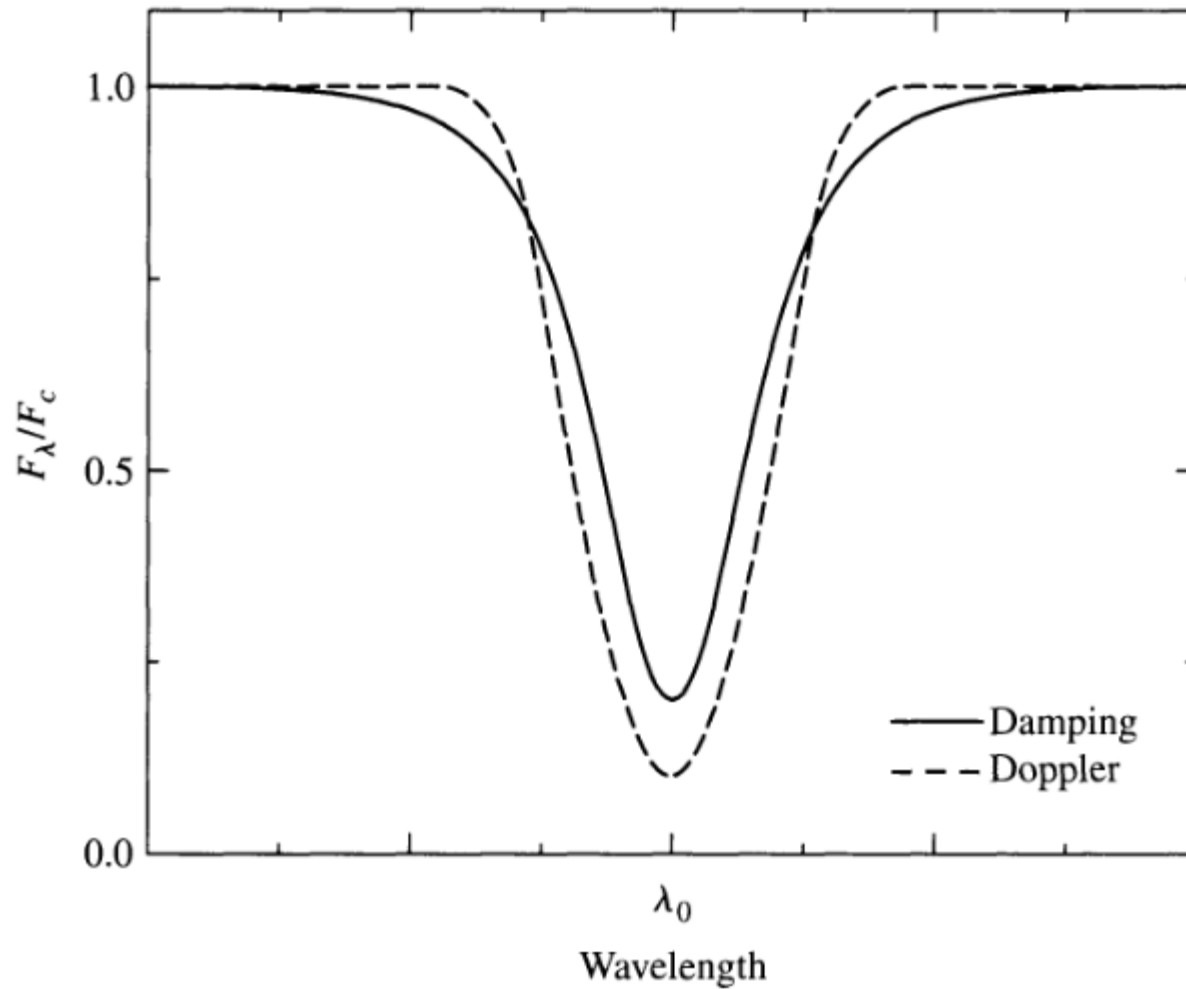
- From thermal motion of atoms
- $\sim 4.27 \times 10^{-2}$  nm at 5772K

- **Pressure and Collisional Broadening**

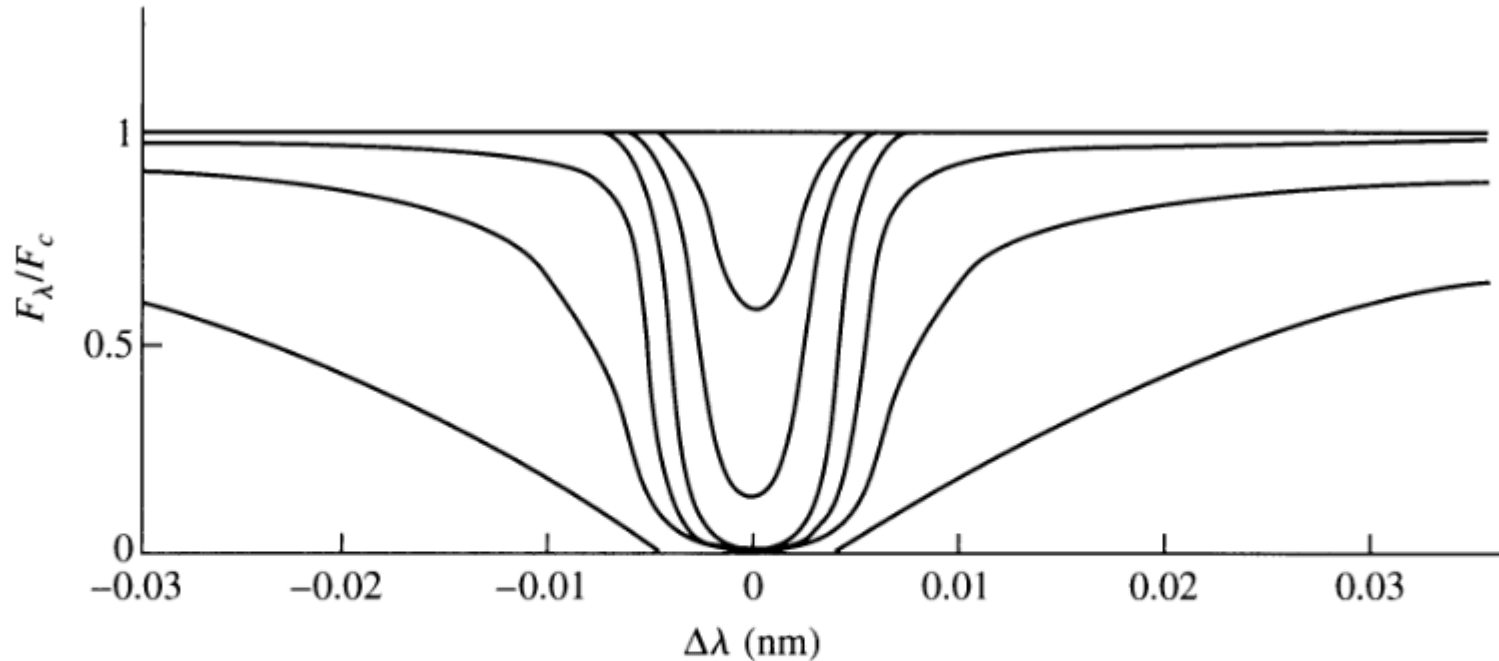
- From atomic orbitals being perturbed from collisions
- $\sim 2-4 \times 10^{-5}$  nm

# The Voigt Profile

- Combination of all broadening effects



# The Voigt Profile



**FIGURE 9.20** Voigt profiles of the K line of Ca II. The shallowest line is produced by  $N_a = 3.4 \times 10^{15}$  ions  $\text{m}^{-2}$ , and the ions are ten times more abundant for each successively broader line. (Adapted from Novotny, *Introduction to Stellar Atmospheres and Interiors*, Oxford University Press, New York, 1973.)